

**Citation:** Wang L, Cai L, Xiao H. A novel method to obtain Manson–Coffin law of metallic materials for low-cycle fatigue by small punch test. *Journal of Harbin Institute of Technology (New Series)*. DOI: 10.11916/j.issn.1005–9113.25039

# A Novel Method to Obtain Manson-Coffin Law of Metallic Materials for Low-Cycle Fatigue by Small Punch Test

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**Abstract:** The assessment of materials low-cycle fatigue (LCF) properties using micro-samples extracted from in-service structures is crucial for structural integrity evaluation. Existing small punch tests for low cycle fatigue (SPT<sub>LCF</sub>) methodologies predominantly rely on load-controlled approaches and remain constrained by dependencies on conventional standard round bar specimen data. To address this limitation, we propose a displacement-controlled (strain-controlled) SPT<sub>LCF</sub> framework. Firstly, a disc with center hole (Disc<sub>ch</sub>) specimen with dual mechanically fixed boundaries (inner and outer circular boundaries) is developed to enable symmetrical displacement-controlled cyclic loading. Then, a comprehensive theoretical framework is presented, including displacement-load model, equivalent stress-strain model, stress-strain field model, and maximum stress-strain model. Additionally, a new approach is proposed to determine material stress-strain relationships under cyclic steady states, coupled with a novel SPT<sub>LCF</sub> methodology for LCF performance assessment. Finite element analysis (FEA) of the specified materials demonstrates less than 4% deviation from model predictions. Experiment results of 316L stainless steel and P91 steel validate that the cyclic steady-state stress-strain relationships and the Manson-Coffin law derived from Disc<sub>ch</sub> specimens align closely with those from standard round bar. This study provides a reliable and material-efficient approach to evaluate LCF properties in metallic materials using miniaturized specimens, particularly for in-service structural applications.

**Keywords:** SPT; Manson-Coffin law; LCF; theoretical model; stress-strain relation

**CLC number:** TG115

**Document code:** A

**Article ID:** 1005-9113(2025)00-0000-17

## 0 Introduction

In critical engineering domains such as aviation and nuclear power, due to the constraints imposed by the dimensions of service structural materials, the method of obtaining low cycle fatigue performance of materials using standard round bar specimens (measured in centimeters and decimeters)<sup>[1]</sup> is difficult to apply. Consequently, small sample fatigue testing technology for structural miniaturization has become a common trend.

Early studies indicated that straight thin sheet specimens are prone to cyclic instability when strain amplitudes exceed 0.2%–0.5%. To enhance specimen stability, Martin<sup>[2]</sup> proposed a constant-radius, funnel-shaped specimen design. Through geometric optimization, which equalized funnel curvature radius

and root width, they achieved a nominal strain instability threshold of 1.5%. Wisner et al.<sup>[3]</sup> validated this design using zirconium alloy sheets, demonstrating effective instability suppression at strain amplitude below 1.41%. However, their life prediction method, which was based on mean strain and stress amplitude, contained a critical limitation; it neglected the stress concentration at the funnel root. This oversight led to deviation from the Manson-Coffin law derived from standard cylindrical specimens.

Addressing the equivalence challenge between thin sheet specimens and standard specimens, Huang et al.<sup>[4-5]</sup> proposed a lateral span displacement-controlled test method. Using finite element analysis (FEA), they derived equations to convert axial strain to nominal strain, achieving consistent low-cycle fatigue data between sheet and cylindrical specimens. However, the phenomenon of material cyclic

Received 2025–08–01.

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hardening and softening during the cyclic process was not considered in the finite element analysis. Jia et al.<sup>[6]</sup> further minimized discrepancies with standard specimen results by incorporating cyclic stress-strain relationships in FEA. Yin et al.<sup>[7]</sup> introduced a semi-analytical method based on energy separation. This technique allowed them to derive the relationship between cyclic stress and strain directly from the relationship between load and displacement. This breakthrough enabled the prediction of low-cycle fatigue life for SS316L stainless steel by utilizing strain amplitudes of representative volume elements (RVE) at funnel roots. Theoretical validation was achieved through the energy density equivalence principles established by Chen and Cai<sup>[8-9]</sup>. In the same year, Liu et al.<sup>[10-11]</sup> developed a theoretical framework for Martin-type specimens based on energy density equivalence and finite element analysis (FEA), characterizing the relationships among cyclic constitutive material parameters, geometric dimensions, displacement amplitude, and load amplitude. However, their theoretical model exhibited limited universality in terms of materials and geometry. Building on energy density equivalence and dimensional analysis, Zhang et al.<sup>[12-13]</sup> formulated analytical expressions for equivalent stress and equivalent strain at the median point of RVE energy density, represented by lateral displacement and load. They introduced a novel low-cycle fatigue testing approach using thin funnel specimens, which demonstrated broader applicability.

SPT<sub>LCF</sub> technology originated from experimental exploration of non-metallic materials. Villarraga et al.<sup>[14-15]</sup> were the first to perform load-controlled SPT<sub>LCF</sub> on polyethylene disks, establishing correlations among fatigue life, peak load, and hysteresis energy. Xiong et al.<sup>[16]</sup> analyzed the influence of experimental factors on fatigue behavior using alumina ceramic square specimens. With industrial demands, SPT<sub>LCF</sub> research shifted to metallic materials. Prakash et al.<sup>[17]</sup> used FEA to identify that the maximum stress under cyclic loading concentrates within a region one-fifth of the punch diameter in disk specimens. They proposed corrections for simulated fatigue life and defined failure based on surges in hysteresis energy. Lancaster et al.<sup>[18]</sup> designed a dual-punch system for symmetric cyclic loading. Subsequent studies on disk specimens under various loading ratios demonstrated consistent damage mechanisms between SPT<sub>LCF</sub> and uniaxial

fatigue tests. By integrating FEA with empirical correlation methods, relationships between maximum principal stress and fatigue life were established at a stress ratio of 0.1. Nonetheless, the empirical basis of these correlations limited their theoretical robustness<sup>[19-20]</sup>. Komazaki et al.<sup>[21-22]</sup> employed a hydraulic test chamber to apply symmetric stress cycling on indented circular specimens, defining failure criteria as either a drop in hydraulic pressure or through-thickness cracking. Due to challenges in measuring critical strain, using average strain values for life prediction introduced bias.

Recent research has focused on developing life prediction frameworks for SPT<sub>LCF</sub>. Zhao et al.<sup>[23]</sup> conducted high-temperature load cycling tests on circular specimens with a loading ratio of 0.1. By substituting stress-related parameters with corresponding load parameters, they evaluated the applicability of conventional life prediction models to SPT<sub>LCF</sub> data. In the same year, Lu<sup>[24]</sup> proposed a critical-plane energy parameter method for load-controlled SPT<sub>LCF</sub> at a loading ratio of 0.1, achieving relatively accurate, albeit conservative, predictions. Otero et al.<sup>[25]</sup> developed an SPT<sub>LCF</sub> methodology using square specimens, correlating compliance increments with damage progression under load-controlled conditions with a loading ratio of 0.1. They introduced a correlation coefficient  $k_{SPT}$  between SPT<sub>LCF</sub> and rotating beam fatigue limits, accounting for specimen dimensions and stress distributions. Kim et al.<sup>[26]</sup> performed symmetric load-controlled SPT<sub>LCF</sub> simulations on square sheets using a modified Wahab damage model. By converting maximum loads into equivalent stresses, they derived S-N curves and validated the reliability of SPT<sub>LCF</sub> through comparison with uniaxial data. Pandey et al.<sup>[27]</sup> systematically examined the effects of load and stress parameters on ratcheting behavior in load-controlled SPT<sub>LCF</sub> under a stress ratio of 0.1. Leveraging three established life prediction models—Gerber, Goodman, and Smith-Watson-Topper—they proposed new damage parameters to correlate ratcheting effects in SPT<sub>LCF</sub> with uniaxial fatigue test results.

Currently, most SPT<sub>LCF</sub> methods focus on load-controlled conditions, with life predictions largely dependent on uniaxial data and lacking rigorous theoretical foundations. Moreover, the implementation of strain-controlled SPT<sub>LCF</sub> still faces critical limitations.

To address the research gap in strain-controlled

small punch testing for low-cycle fatigue, this study designed a disc specimen with a center hole ( $Disc_{ch}$ ) capable of displacement-controlled (i. e., strain-controlled) loading, along with a corresponding clamping system. Based on the energy density equivalence principle and validated by FEA, a theoretical framework was established comprising: (1) theoretical models for load-displacement and stress-strain relationships, and (2) a semi-analytical model characterizing stress-strain fields and maximum stress-strain values. The framework was calibrated and validated through FEA. Furthermore, novel methods were proposed for obtaining material cyclic stress-strain relationships and implementing strain-controlled small punch fatigue testing. The reliability of these methods was confirmed via comparative analysis between cyclic small punch tests on P91 and 316L materials and standard uniaxial test results.

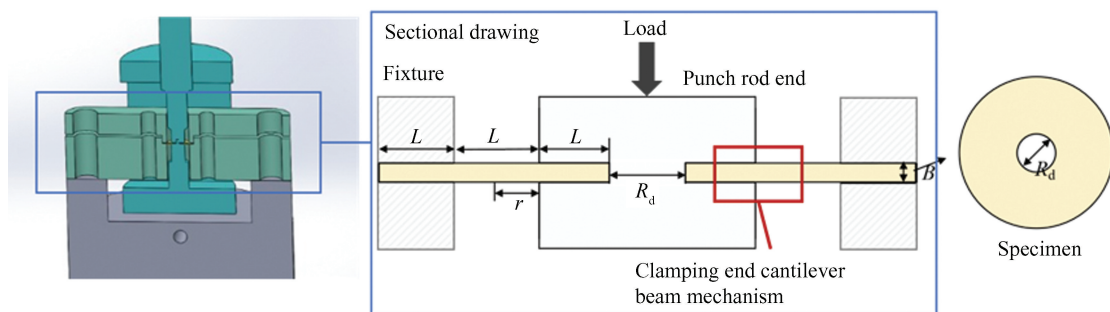
## 1 Experiment Setup

### 1.1 Finite Element Analysis Conditions

The schematic illustration in Fig.1 depicts the  $Disc_{ch}$  specimen and the associated  $SPT_{LCF}$  test setup. The critical geometrical parameters are defined as follows:  $R_d$  denotes the inner diameter of the central aperture in the annular specimen,  $r$  represents the

radial distance measured from the punch clamping interface,  $B$  corresponds to the axial thickness of the specimen, and  $L$  specifies the effective clamping length between the punch termination point and the specimen-gripping region. It should be noted that both upper and lower fixtures maintain identical punching configurations, ensuring symmetrical loading conditions during experimental operations.

As highlighted in the red box in Fig.1, the clamping end behaves in analogous to a cantilever beam. Stress analysis reveals that the supported regions on the upper and lower surfaces of the fixed end are dominated by normal stress: the upper surface is subjected to maximum tensile stress, while the lower surface experiences maximum compressive stress, both of equal magnitude. Near the fixed end, the shear stress component becomes negligible, rendering the von Mises equivalent stress equal in magnitude to the principal stress. This uniaxial stress state provides a theoretical basis for employing the local stress-strain method to determine the equivalent stress and strain amplitudes under cyclic steady-state conditions. Moreover, during cyclic loading, the specimen end in contact with the punch undergoes uniaxial proportional loading, and the stress evolution at the critical failure point exhibits characteristic features of uniaxial loading.



**Fig.1 Schematic representation of the  $Disc_{ch}$  specimen and  $SPT_{LCF}$  fixture system**

FEA was conducted in ANSYS 19.2 to simulate the mechanical response of  $Disc_{ch}$  specimens under monotonic loading. The material was modeled as homogeneous, isotropic, and following an isotropic hardening rule, with a Poisson's ratio of 0.3. The mesh was discretized using PLANE 182 elements, selected for their suitability in axisymmetric simulations. To accurately resolve stress concentrations and deformation gradients, local mesh refinement was applied in the contact regions between the specimen

and the punch. The punch rod was modeled as a rigid body with TARGE 169 elements, and the contact interfaces between the  $Disc_{ch}$  specimen and the upper and lower fixtures were represented using CONTA 172 flexible contact elements, forming contact pairs to simulate mechanical interaction. Boundary conditions were assigned to replicate experimental constraints: axisymmetric conditions were applied to the left and bottom edges of the model, while the upper and lower fixtures were fully fixed ( $X = 0, Y = 0$ ) to mimic

clamping. A friction coefficient of 0.15 was defined at the punch-specimen contact interface to account for

surface interactions. The mesh configuration result is depicted in Fig.2.

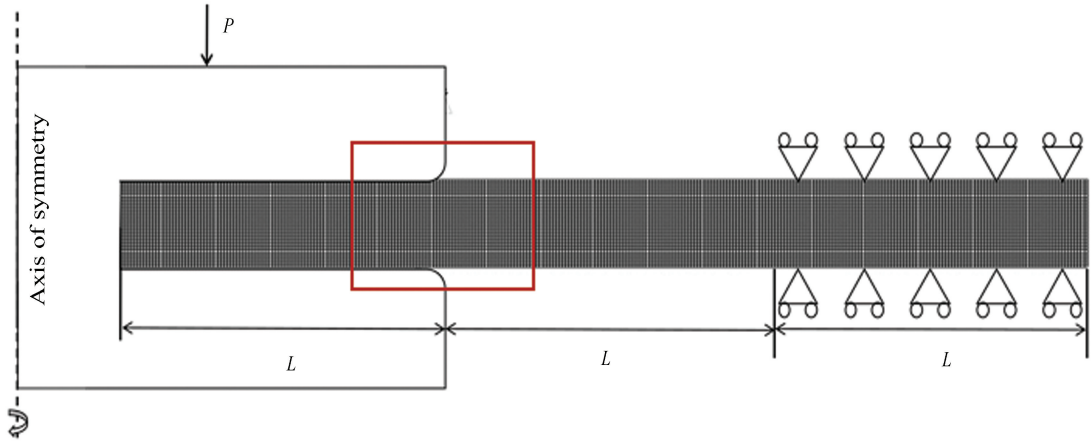
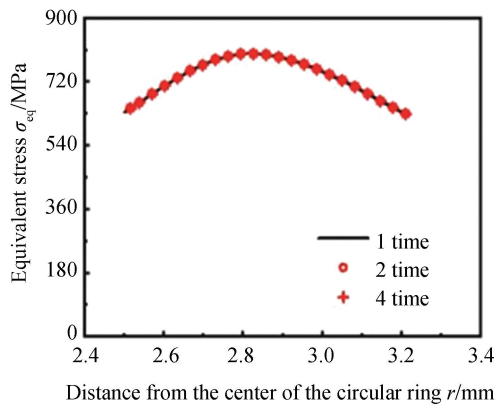
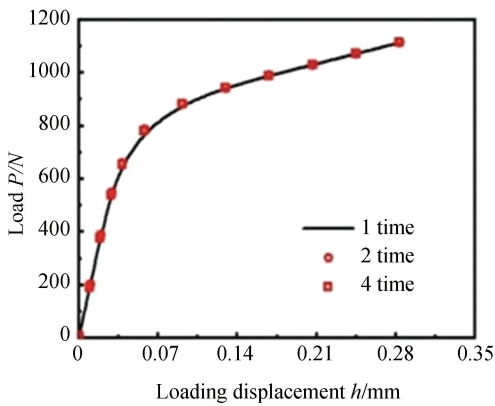


Fig.2 Disc<sub>ch</sub> specimen symmetric mesh model

To investigate the influence of mesh density on FEA results and to establish an optimal mesh configuration that ensures computational convergence, a comprehensive finite element analysis was performed on H-law materials with the following mechanical properties: Young's modulus  $E = 200$  GPa, Hardening index  $n = 0.2$ , and yield strength  $\sigma_y = 400$  MPa. Fig. 3 presents the load-displacement

responses and equivalent stress distributions under different mesh densities, showing converged results for densities of 1, 2, and 4. To optimize computational efficiency while maintaining solution accuracy and enabling high-resolution contour plot analysis, a mesh density of 2 was selected for detailed analysis, comprising a total of 36384 elements.



(a) Load-displacement response under different grid densities (b) Equivalent stress distribution under different grid densities

Fig.3 Influence of grid density on FEA results

## 1.2 Test Conditions

The geometric dimensions of sample and punch end are illustrated in Fig. 4. As shown in Fig. 4(a), the central-hole disc specimen has an inner diameter of 2 mm, an outer diameter of 14 mm, and a thickness of 0.5 mm. Fig.4(b) presents the dimensional parameters of the punch end configuration. The contact surfaces between the specimen and both the upper and lower punches have a diameter of 5.6 mm. The edges of the contact

regions are designed with a fillet radius of 0.5 mm and a transition zone width of 0.2 mm.

Fig.5 illustrates the clamping system for the central-hole disc specimen. The system primarily comprises two assemblies: 1) a fixed strip set, shown in Fig. 5 (a), which clamps the outer circumferential edge of the specimen; and 2) upper and lower punch rods combined with a loading connector, shown in Fig. 5 (b), which clamp the inner hole edge of the specimen.

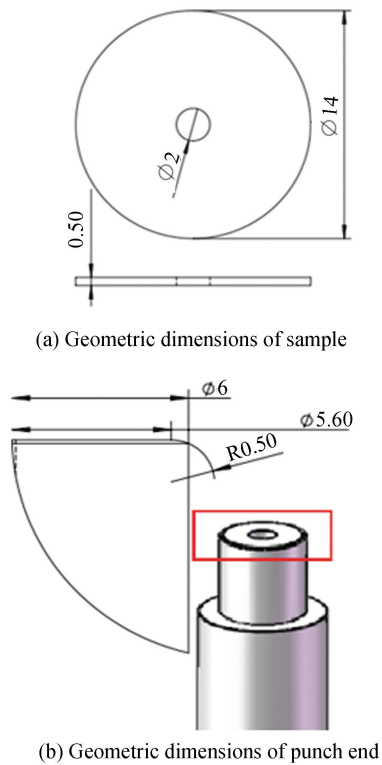


Fig.4 Geometric dimensions of sample and punch end

The outer circumferential edge of the specimen 5 is secured within a sample cavity between the upper 3 and lower 4 fixed strips, with a radial clamping width of 2 mm. The cavity has a diameter of 14 mm and a height of 3 mm. The assembled fixed strip pair is inserted into the loading connector, while the upper 1 and lower 2 punch rods pass through it. The opposing ends of the punch rods contact the upper and lower surfaces of the specimen, respectively, clamping the inner hole periphery with a radial width of 2 mm.

The upper and lower punch rods are fastened via the loading connector 6. The fixed strip pair is bolted to the lower clamp of the testing machine 7, and the upper punch rod is connected to the upper clamp 8 to allow vertical movement. Relative displacement between the punch rods and the fixed strip pair introduces controlled displacement at both the inner hole periphery and the outer circumferential clamping regions. This setup enables displacement-controlled cyclic loading of the central-hole disc specimen, with the relative displacement measured by a strain extensometer 9.

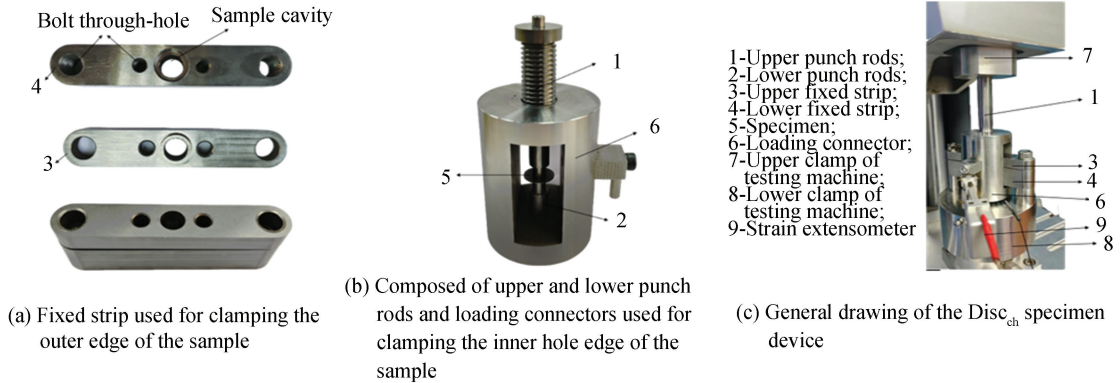


Fig. 5 Disc<sub>ch</sub> specimen clamping system

To validate the uniqueness of stress-strain relationship for metallic materials under cyclically stable conditions and to assess the applicability of the novel SPT<sub>LCF</sub> method, a series of multi-stage symmetric cyclic displacement-controlled small punch tests (SPT) were performed on Disc<sub>ch</sub> specimens. The experimental results were compared to the results of corresponding standard round bar specimens. The SPT<sub>LCF</sub> were conducted using a Gouyuan testing system, manufactured by Chengdu Miniature Mechanical Testing Science & Technology Co., Ltd. (MMTS). The load sensor of the equipment has a capacity of  $\pm 5$  kN. Displacement control was achieved using an MTS32.29F-30 strain extensometer

with a gauge length of 5 mm and a range of  $-10\%$  to  $30\%$ . All tests were performed under a loading ratio of  $R = -1$ .

## 2 Theoretical Model

### 2.1 Load-Displacement Model

The elastoplastic stress-strain relationship for commonly used engineering metals is frequently modeled by the H-law model, expressed as Eq.(1).

$$\sigma_{eq} = \begin{cases} E \varepsilon_{eq} & \varepsilon_{eq} < \varepsilon_y \\ K \varepsilon_{eq}^n & \varepsilon_{eq} \geq \varepsilon_y \end{cases} \quad (1)$$

where,  $\varepsilon_{eq}$  and  $\sigma_{eq}$  represent the equivalent strain and stress, respectively.  $\varepsilon_y$  is yield strain;  $E$  denotes the

elastic modulus,  $n$  is the strain hardening index, and  $K$ , the strain hardening coefficient, is defined as  $K = E^n \sigma_y^{1-n}$ .

For a homogeneous and isotropic solid under uniaxial loading, the relationship between median energy density and total deformation energy is given by the following Eq.(2) [28-29],

$$U = V_{\text{eff}} u_C = V_{\text{eff}} \int_0^{\varepsilon_{\text{eq-C}}} \sigma_{\text{eq-C}} d\varepsilon_{\text{eq-C}} \quad (2)$$

here,  $V_{\text{eff}}$  is the effective deformation volume,  $u_C$  the energy density at the midpoint C of the the representative volume element (RVE<sub>C</sub>),  $\sigma_{\text{eq-C}}$  the equivalent stress of RVE<sub>C</sub> under specific loading conditions, and  $\varepsilon_{\text{eq-C}}$  the equivalent strain of RVE<sub>C</sub>.

Combined with the H-law Eq.(1) and Eq.(2), the strain energy density of RVE<sub>C</sub> is calculated as Eq.(3):

$$UC_C = \int_0^{\varepsilon_y} \sigma_{\text{eq-C}} d\varepsilon_{\text{eq-C}} + \int_{\varepsilon_y}^{\varepsilon_{\text{eq-C}}} \sigma_{\text{eq-C}} d\varepsilon_{\text{eq-C}} = \frac{E^n \sigma_y^{1-n} \varepsilon_{\text{eq-C}}^{1+n}}{1+n} \left[ 1 - \frac{1-n}{2} \left( \frac{\varepsilon_y}{\varepsilon_{\text{eq-C}}} \right)^{1+n} \right] \quad (3)$$

Under large deformation conditions,  $\varepsilon_y$  is relatively small compared to  $\varepsilon_{\text{eq-C}}$ , as demonstrated in Eq.(4).

$$\frac{1-n}{2} \left( \frac{\varepsilon_y}{\varepsilon_{\text{eq-C}}} \right)^{1+n} \ll 1 \quad (4)$$

Simplifies Eq.(3) into Eq.(5).

$$u_C = \frac{K \varepsilon_{\text{eq-C}}^{1+n}}{1+n} = \frac{\sigma_{\text{eq-C}} \varepsilon_{\text{eq-C}}}{1+n} \quad (5)$$

Under the application of load P, when a solid specimen undergoes a displacement  $h^*$  along the loading axis at the loaded position, both  $V^*$  and  $\varepsilon_{\text{eq-C}}$  exhibit co-varying adjustments. The characteristic volume  $V^*$ , characteristic area  $A^*$ , and characteristic geometric dimension  $h^*$  are defined to satisfy the relationship:  $V^* = A^* h^*$  and it is hypothesized that both  $\varepsilon_{\text{eq-C}}$  and  $V_{\text{eff}}/V^*$  follow power-law relationships with  $h/h^*$ :

$$\begin{cases} \frac{V_{\text{eff}}}{V^*} = k_1 \left( \frac{h}{h^*} \right)^{k_2} \\ \varepsilon_{\text{eq-C}} = k_3 \left( \frac{h}{h^*} \right)^{k_4} \\ k_1 = p_1 + p_2 \varepsilon_y \end{cases} \quad (6)$$

where  $k_1$  and  $k_2$  are the effective deformation volume coefficient and index respectively,  $k_3$  and  $k_4$  are the equivalent deformation coefficient and index respectively,  $p_1$  and  $p_2$  are effective deformation

volume correction parameter.

Under quasi-static loading conditions, energy conservation implies that the work  $W$  done by the external load on the specimen equals the total deformation energy  $U$ ,

$$\begin{cases} W = U \\ \int_0^h P dh = V_{\text{eff}} u_C \end{cases} \quad (7)$$

According to Eq. (6), both  $V_{\text{eff}}$  and  $u_C$  are functions of the  $V$  and  $h$ . Through application of the differentiation operator with respect to  $h$  on both members of Eq. (7),  $P$  can be described as in Eq.(8).

$$P = u_C \frac{dV_{\text{eff}}}{dh} + V_{\text{eff}} \frac{du_C}{dh} \quad (8)$$

We solve Eq.(6) in conjunction with Eqs.(5) and (1), resulting in the following integral expression for  $u_C$ ,

$$u_C = \int_0^h \frac{k_3 k_4}{h^*} \left( \frac{h}{h^*} \right) \sigma_{\text{eq-C}} dh \quad (9)$$

Substituting  $u_C$  from Eq. (5) and  $V_{\text{eff}}$  from Eq.(6) into the first term on the right-hand side of Eq.(8), and similarly substituting  $u_C$  from Eq.(9) and  $V_{\text{eff}}$  from Eq.(6) into the second term, yields,

$$P = \sigma_{\text{eq-C}} A^* k_1 k_3 \left( \frac{k_2}{n+1} + k_4 \right) \left( \frac{h}{h^*} \right)^{k_2+k_4-1} \quad (10)$$

Rearranging the above formula yields Eq. (11).

$$\begin{cases} \frac{P}{P^*} = C \left( \frac{h}{h^*} \right)^m \\ P^* = KA^* \\ C = \frac{k_1 k_3^{n+1} (m+1)}{n+1} \\ m = k_4 (n+1) + k_2 - 1 \\ k_1 = p_1 + p_2 \varepsilon_y \end{cases} \quad (11)$$

where,  $C$  and  $m$  serve as the loading coefficient and index, respectively, for the second stage of the load-displacement curve during SPT<sub>LCF</sub> for the Disc<sub>ch</sub> model, and  $P^*$  represents the characteristic load. This equation encapsulates the relationship between the H-law parameters, geometric dimensions, and the load and displacement of the Disc<sub>ch</sub> specimen, called as load-displacement model.

## 2.2 Equivalent Stress and Strain Model

Utilizing Eqs. (10) and (6), equivalent stress and strain for the RVE<sub>C</sub> are derived as,

$$\left\{ \begin{array}{l} \sigma_{\text{eq-C}} = k_5 \frac{P}{A^*} \left( \frac{h}{h^*} \right)^{k_6} \\ \varepsilon_{\text{eq-C}} = k_3 \left( \frac{h}{h^*} \right)^{k_4} \\ \left[ \begin{array}{l} k_5 = \left[ k_1 k_3 \left( k_4 + \frac{k_2}{n+1} \right) \right]^{-1} \\ k_6 = 1 - k_2 - k_4 \\ k_1 = p_1 + p_2 \varepsilon_y \end{array} \right] \end{array} \right. \quad (12)$$

where,  $k_5$  and  $k_6$  are the equivalent stress coefficient and index respectively. Eq. (12) delineates the relationship among equivalent stress and strain of  $\text{RVE}_C$ , the load, displacement, geometric dimensions, and material hardening index of the specimen, designated as the equivalent stress and strain model.

### 2.3 Experimental Method for Obtaining H Law Parameters of $\text{Disc}_{\text{ch}}$

The loading index  $m$  and coefficient  $C$  are determined through a regression analysis of the second stage loading-displacement curve of the  $\text{Disc}_{\text{ch}}$  specimen in  $\text{SPT}_{\text{LCF}}$ , expressed as  $P/P^* = C (h/h^*)^m$ . From Eqs. (11) and (1) we derive the solution for the H-law parameters  $n$ ,  $K$ , and  $\sigma_y$ .

$$\left\{ \begin{array}{l} n = \frac{m+1-k_2}{k_4} - 1 \\ K = \frac{P^*}{A^*} \\ \sigma_y = \left( \frac{P^*}{E^n A^*} \right)^{\frac{1}{1-n}} \end{array} \right. \quad (13)$$

### 2.4 Stress and Strain Fields Model

Using dimensional analysis, the stress and strain distributions near the punch's clamping end are described by Eq. (14) below, with  $\sigma_{\text{eq-C}}$  and  $\varepsilon_{\text{eq-C}}$  representing the characteristic stress ( $\sigma_C$ ) and strain ( $\varepsilon_C$ ).

$$\left\{ \begin{array}{l} \frac{\sigma_{\text{eq}}}{\sigma_C} = f(r/R_d, n, \varepsilon_y) \\ \frac{\varepsilon_{\text{eq}}}{\varepsilon_C} = (r/R_d, n, \varepsilon_y) \end{array} \right. \quad (14)$$

where,  $\sigma_C$  and  $\varepsilon_C$  contain the load  $P$ , loading displacement  $h$ , hardening index  $n$ ,  $g (r/R_d, n, \sigma_y/E)$  and  $f (r/R_d, n, \sigma_y/E)$  contain the distance  $r$  from the punch clamping endpoint and material parameters (elastic modulus  $E$ , strain hardening index  $n$ , yield stress  $\sigma_y$ ).  $R_d$  is the inner diameter of the mesoporous circular specimen.

To effectively characterize the distinctive

distribution characteristics of the stress and strain fields in  $\text{Disc}_{\text{ch}}$  specimens under cyclic loading, a trigonometric function form is adopted in this study to construct the constitutive relation. Compared to other forms, such as polynomials and exponentials, the trigonometric function offers higher parameter sensitivity and better operational tolerance when fitting field distributions. It allows flexible adaptation to variations in load response under different materials and displacement conditions through adjustments of parameters such as amplitude and frequency. Moreover, the incorporation of a constant term enables the representation of non-zero mean or notable bias characteristics in the distribution of stress and strain fields, thereby providing a more accurate description of their physical nature. Functions  $f$  and  $g$  are assumed as follow:

$$\left\{ \begin{array}{l} f = \cos \left[ a_1 \left( \frac{r}{R_d} \right) + a_2 \right] + a_3 \\ g = c_1 \cos \left[ c_2 \left( \frac{r}{R_d} \right) + c_3 \right] \end{array} \right. \quad (15)$$

where  $a_k, (k=1,2,3)$  and  $c_s, (s=1,2,3)$  represent the dimensionless parameters related to material parameters. Eqs. (14) and (15) elucidate the relationship between the equivalent stress and strain at different positions within the  $\text{Disc}_{\text{ch}}$  specimen and its associated load, displacement, geometric dimensions, material parameters, and distance from the punch clamping endpoint, known as stress and strain fields model.

### 2.5 Maximum Stress and Strain Model

The sample is clamped between the upper and lower punch clamps for loading. As highlighted in the red box in Fig.2, the mechanism at the clamping end functions similarly to a cantilever beam. Mechanical analysis and FEA result indicate that the stress and strain values peak at the clamping end. By setting  $r/R_d=0$  in the stress and strain field equations (Eqs. (14) and (15)), the maximum stress  $\sigma_{\text{eq-max}}$  and maximum strain  $\varepsilon_{\text{eq-max}}$  are obtained, as shown in Eq.(16).

$$\left\{ \begin{array}{l} \frac{\sigma_{\text{eq-max}}}{\sigma_C} = 1 + a_3 \\ \frac{\varepsilon_{\text{eq-max}}}{\varepsilon_C} = c_1 \end{array} \right. \quad (16)$$

Under the applied load  $P$ , integrating Eqs.(12) and (16) establishes the following relationships:

$$\begin{cases} \sigma_{\text{eq-max}} = \beta_1 \frac{P}{A^*} \left( \frac{h}{h^*} \right)^{k_6} \\ \varepsilon_{\text{eq-max}} = \beta_2 \left( \frac{h}{h^*} \right)^{k_4} \\ \left[ k_5 = \left[ k_1 k_3 \left( k_4 + \frac{k_2}{n+1} \right) \right]^{-1} \right. \\ \left. k_6 = 1 - k_2 - k_4 \right. \\ \left. \beta_1 = (1 + a_3) k_5 \right. \\ \left. \beta_2 = c_1 k_3 \right. \end{cases} \quad (17)$$

where  $\beta_1$  and  $\beta_2$  are the maximum stress and maximum strain coefficient, respectively. Eq.(17) clarifies the relationship between the maximum stress, maximum strain of the Disc<sub>ch</sub> specimen, and the loading displacement, load, and material parameters, referred to as the maximum stress and strain model. This relationship is essential for converting SPT<sub>LCF</sub> displacement and load data into maximum stress and strain in the Disc<sub>ch</sub> specimen. In the cyclic loading state, the maximum stress and strain of the specimen represent the stress amplitude and strain amplitude at the failure point. This model facilitates the conversion of the load amplitude  $P_a$  and the axial displacement amplitude  $h_a$  at the center of the sample to the stress amplitude  $\sigma_{\text{eq-a}}$  and strain amplitude  $\varepsilon_{\text{eq-a}}$  at the failure point under low cycle loading.

## 2.6 Equivalent Stress-Equivalent Strain Relationship of Cyclic Steady State

In the steady-state strain cycle, for metal materials that adhere to the Masing effect, the cyclic steady stress amplitude  $\sigma_a$  and strain amplitude  $\varepsilon_a$  can represent the uniaxial stress-strain behavior of materials under cyclic stability. For metals that do not conform to the Masing effect,  $\sigma_a - \varepsilon_a$  can still be used as an approximate measure. Variables such as the load amplitude  $P_a$ , the axial displacement amplitude  $h_a$  at the center of the specimen, the cyclic equivalent stress amplitude  $\sigma_{\text{eq-C-a}}$ , and the cyclic equivalent strain amplitude  $\varepsilon_{\text{eq-C-a}}$ , replace the respective variables  $P$ ,  $h$ ,  $\sigma_{\text{eq-C}}$ ,  $\varepsilon_{\text{eq-C}}$  in Eq. (12). Theoretical equations for the cyclic equivalent stress amplitude and equivalent strain amplitude for the RVE<sub>C</sub>, represented by the axial displacement amplitude at the center of the specimen and the load amplitude under elastic-plastic conditions, can be derived.

$$\begin{cases} \sigma_{\text{eq-C-a}} = k_5 \frac{P_a}{A^*} \left( \frac{h_a}{h^*} \right)^{k_6} \\ \varepsilon_{\text{eq-C-a}} = k_3 \left( \frac{h_a}{h^*} \right)^{k_4} \\ \left[ k_5 = \left[ k_1 k_3 \left( k_4 + \frac{k_2}{n+1} \right) \right]^{-1} \right. \\ \left. k_6 = 1 - k_2 - k_4 \right. \end{cases} \quad (18)$$

Using the symmetric strain cyclic test under different displacement amplitudes as examples, the stress amplitude-strain amplitude relationship for materials in steady-state cyclic conditions can be directly obtained through the axial displacement amplitude at the center of the specimen and the load amplitude. This method involves collecting and processing data, where the center axial displacement-load hysteresis loop of the sample for each controlled displacement amplitude level is recorded. Consequently, the load amplitude and the center axial displacement amplitude data ( $P_a$ ,  $h_a$ ) corresponding to each displacement amplitude level are obtained. The cyclic equivalent stress amplitude and strain amplitude under elastic-plastic conditions are calculated using Eq. (18). Additionally, the expressions for the cyclic strain strengthening index  $n'$  and the cyclic strain strengthening coefficient  $K'$  of the cyclic stress-strain relationship can be derived from Eq. (12). Ultimately, from the theoretical equations of cyclic equivalent stress amplitude and equivalent strain amplitude, the data ( $\sigma_{\text{eq-C-a}}$ ,  $\varepsilon_{\text{eq-C-a}}$ ) for the RVE<sub>C</sub> can be obtained, presenting the relationship between equivalent stress amplitude and equivalent strain amplitude when the material is cyclically stable.

## 3 Parameters Calibration

### 3.1 Parameters Calibration of Load-Displacement Model

Before calibrating the model parameters, it is essential to select appropriate characteristic areas  $A^*$  and characteristic displacements  $h^*$  to mitigate the impact of scale effects on the results. This overlap of the normalized responses,  $P/P^* - h/h^*$  is achieved when testing geometrically similar specimens whose dimensions satisfy the ratio  $L = 4B$ . FEA has established the optimal characteristic parameters for the Disc<sub>ch</sub> specimen as a characteristic displacement of  $h^* = B$  and a characteristic area of  $A^* = 3BL$ .

The elastic modulus  $E$  of the material used for the Disc<sub>ch</sub> specimen was set at 200 GPa. The yield stress  $\sigma_y$  ranged from 200 MPa to 1000 MPa, with increments of 200 MPa, and the strain hardening index  $n$  varied from 0.1 to 0.4, with a step size of 0.1. Through FEA, load-displacement responses for twenty combinations of constitutive parameters  $\sigma_y$  and  $n$  were acquired. Subsequent power law fitting of the second stage establishes relationships between different parameters  $m, C$  and the hardening index  $n$ . According to the load-displacement model of Eq. (11) regression, during the regression analysis, it was observed that the effective deformation volume coefficient  $k_1$  exhibited significant dispersion in materials with varying yield strains  $\varepsilon_y$ . To address this issue, a modified formulation was proposed in Eq.(19) for systematic correction.

$$k_1 = p_1 + p_2 \varepsilon_y \quad (19)$$

By substituting the modified formulation from Eq.(19) into the original theoretical framework of Eq.(11), the complete theoretical model for the load-displacement relationship can be derived as:

$$\begin{cases} \frac{P}{P^*} = C \left( \frac{h}{h^*} \right)^m \\ P^* = KA^* \\ C = \frac{k_1 k_3^{n+1} (m+1)}{n+1} \\ m = k_4 (n+1) + k_2 - 1 \\ k_1 = p_1 + p_2 \varepsilon_y \end{cases} \quad (20)$$

where  $p_1$  and  $p_2$  are the correction parameters of effective deformation volume coefficient. The model parameters  $p_1, p_2, k_3-k_5$  corresponding to the Disc<sub>ch</sub> specimen are obtained, as presented in Table 1.

**Table 1 Parameter values of load-displacement model**

$p_1$	$p_2$	$k_3$	$k_4$	$k_5$
36.7581	0.0457	0.0253	1.1099	4838.1083

It should be noted that during the process of deriving the load-displacement model from Eq. (4) BP)], the approximation of small quantities  $\frac{1-n}{2}$ .

$\left( \frac{\varepsilon_y}{\varepsilon_{eq-C}} \right)^{1+n}$  are neglected so that slight deviations may be exhibited in predicting load-displacement,  $n$ , and  $K$ . Hence, there is a minor discrepancy between the loading index  $m$  calculated directly from Eq.(11) and the loading index  $m_i$  fitted in the second stage of the load-displacement curve. It is recommended to adjust

the loading index by the yield strain  $\varepsilon_y$ ,

$$m_i = p_3 + p_4 \varepsilon_y + p_5 m \quad (21)$$

where  $p_3-p_5$  denote correction parameters, and the corrected parameter values are listed in Table 2.

**Table 2  $m_i$  correction parameters**

$p_3$	$p_4$	$p_5$
-0.0348	-5.2256	5879.8077

### 3.2 Parameter Calibration of Stress-Strain Field Model

Through finite element analysis and mathematical analysis, the relationship between material parameters and  $a_k, c_s$  satisfies the following equation:

$$\begin{cases} a_k = b_{k1} (\varepsilon_y + b_{k2})^{b_{k3n+b_{k4}}} + b_{k5} \\ c_s = d_{s1} (\varepsilon_y + d_{s2})^{d_{s3n+d_{s4}}} + d_{s5} \end{cases} \quad (22)$$

Here,  $b_{kj,j=1-5}, d_{st,t=1-5}$  are constants of the equivalent transformation theoretical model for stress and strain distribution. By combining Eq. (14) and Eq.(15), the complete expression of the equivalent transformation theoretical model for maximum stress (amplitude) and strain (amplitude) can be derived as:

$$\begin{cases} \sigma_{eq-max} = \beta_1 \frac{P}{A^*} \left( \frac{h}{h^*} \right)^{k_6} \\ \varepsilon_{eq-max} = \beta_2 \left( \frac{h}{h^*} \right)^{k_4} \\ \left[ k_5 = \left[ k_1 k_3 \left( k_4 + \frac{k_2}{n+1} \right) \right]^{-1} \right. \\ k_6 = 1 - k_2 - k_4 \\ k_1 = p_1 + p_2 \varepsilon_y \\ \beta_1 = (1 + a_3) k_5 \\ \beta_2 = c_1 k_3 \\ \left. \begin{cases} a_k = b_{k1} (\varepsilon_y + b_{k2})^{b_{k3n+b_{k4}}} + b_{k5} \\ c_s = d_{s1} (\varepsilon_y + d_{s2})^{d_{s3n+d_{s4}}} + d_{s5} \end{cases} \right. \end{cases} \quad (23)$$

The elastic modulus  $E$  of the material was set to 200 GPa, the yield stress  $\sigma_y$  varied from 200 MPa to 1000 MPa in increments of 200 MPa, the strain hardening index  $n$  ranged from 0.1 to 0.4 with a step size of 0.1. Under power-law elastoplastic conditions, FEA was employed to obtain the equivalent stress and equivalent strain distributions of the Disc<sub>ch</sub> specimen for a total of 20 combinations of the constitutive parameters  $\sigma_y$  and  $n$ . Using  $\sigma_C$  and  $\varepsilon_C$  as characteristic quantities, regression was performed based on Eq.(23) to determine the parameter values of  $b_{kj}$  and  $d_{st}$ , which are provided in Table 3 and Table 4, respectively.

**Table 3 Stress field model parameters  $b_{kj}$  values**

$j$	1	2	3	4	5
$b_{1j}$	18.3683	0.9954	-169.7827	88.6527	-11.0653
$b_{2j}$	-0.0055	0.9966	863.9708	-393.7378	-1.3885
$b_{3j}$	-0.4418	0.9995	-2566.7231	-249.3748	0.6839

**Table 4 Strain field model parameter  $d_{st}$  values**

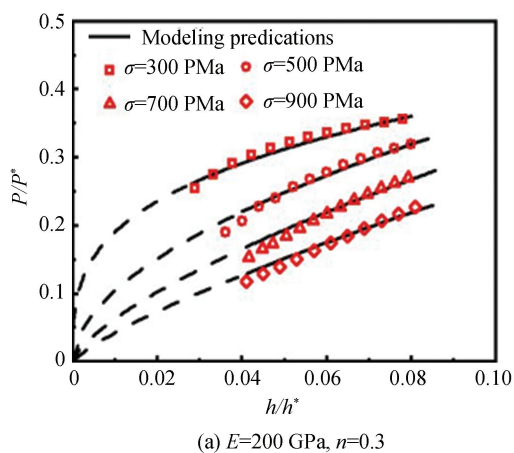
$t$	1	2	3	4	5
$d_{1t}$	0.0377	0.9330	57.3532	-92.2506	2.3669
$d_{2t}$	-5.5355	0.9951	-297.7037	241.9345	11.2167
$d_{3t}$	0.6446	0.9951	-305.8787	284.0280	-0.9371

## 4 Model Validation

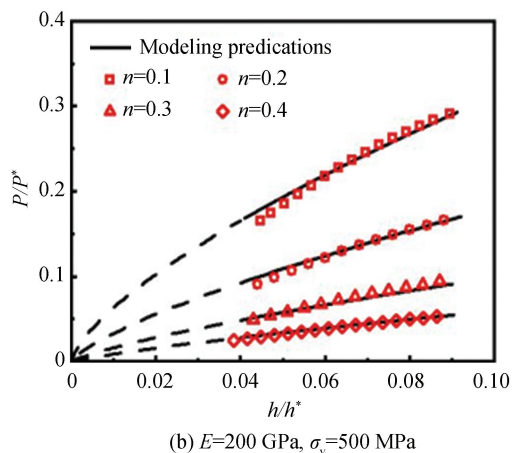
### 4.1 Finite Element Analysis Validation

#### 4.1.1 Verification of load-displacement model

To validate the universality and predictive



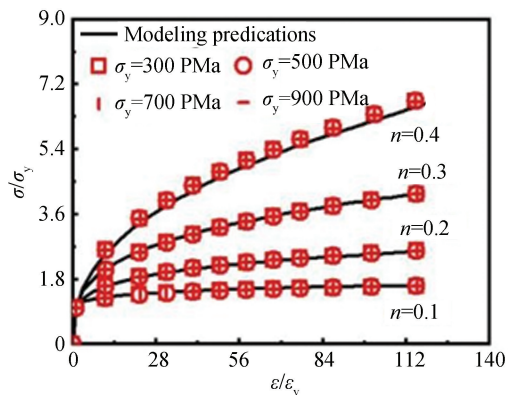
accuracy of the load-displacement model across a range of materials, systematic simulations were conducted with the elastic modulus  $E$  of the sample material set to 200 GPa. The yield stress  $\sigma_y$  was varied from 300 MPa to 900 MPa in increments of 200 MPa, and the was assigned values from 0.1 to 0.4 in steps of 0.1. FEA was performed to obtain the load-displacement response for various combinations of  $\sigma_y$  and  $n$ . The loading coefficient  $C$  and loading index  $m$  for the respective materials are calculated using Eq. (11), which incorporates parameters  $k_1$  to  $k_4$  from Table 1. Based on these values, the  $P$ - $h$  curve predicted by the load-displacement model was plotted. As illustrated in Fig. 6, the model results show excellent agreement with the FEA data, achieving a goodness-of-fit greater than 0.96.



**Fig. 6 Comparison between the predicted results of the load-displacement model and the FEA results under different material parameters ( $E$ ,  $\sigma_y$ ,  $n$ )**

Reverse validation involves fitting a Stage II curve segment of the  $Disc_{ch}$  sample in SPT monotonic loading with the function  $P/P^* = C(h/h^*)^m$ , from which the loading index  $m$  and coefficient  $C$  are determined. According to Eq. (13), the H-law parameters  $n$ ,  $K$ , and  $\sigma_y$  are derived. The outcomes of this inverse analysis are then graphically compared with the known H-law parameters, as displayed in Fig. 7. The results demonstrate that the inverse prediction performs well.

In summary, across a broad spectrum of H-law constitutive parameters typical of conventional engineering materials, both the load-displacement model, and the reverse prediction of H-law parameters for the  $Disc_{ch}$  sample in SPT<sub>LCF</sub> show close agreement with the FEA curves.



**Fig.7 Comparison of reverse predicted values of load-displacement model and true values under different material parameters ( $E$ ,  $\sigma_y$ ,  $n$ )**

#### 4.1.2 Verification of stress-strain field model

The elastic modulus  $E$  of the compressed material

was preset to 200 GPa. The yield stress  $\sigma_y$  was varied from 300 MPa to 900 MPa in increments of 200 MPa, and the strain hardening index  $n$  was assigned values from 0.1 to 0.4 in steps of 0.1. Scatter plots were generated to represent the dimensionless stress  $\sigma_{eq}/\sigma_C$ , dimensionless strain  $\varepsilon_{eq}/\varepsilon_C$ , and radial

position ratio  $r/R$ . Based on Eq.(15), all computational results were compared to the FEA, as depicted in Fig.8 and Fig.9. The results indicate that the model predictions are in close agreement with the FEA results, with a goodness-of-fit exceeding 0.95.

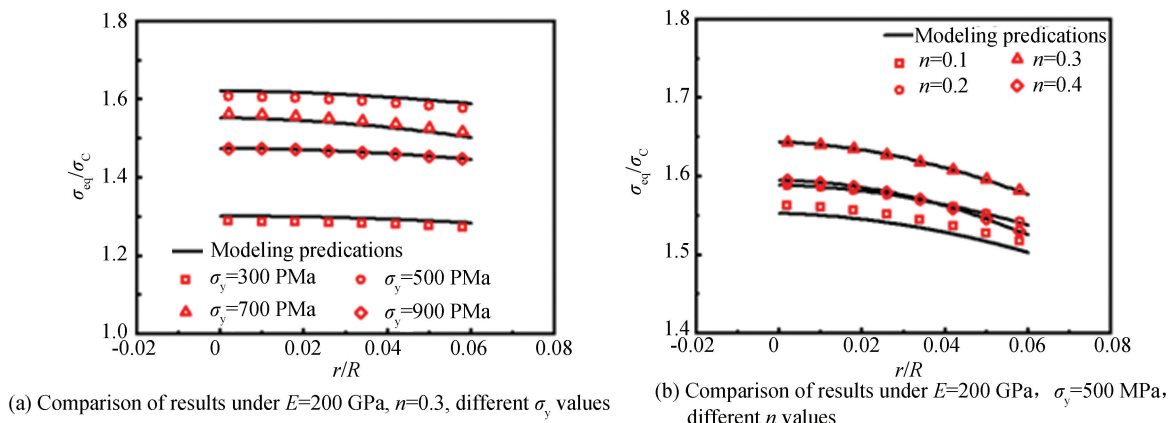


Fig.8 Comparison of stress field model prediction results and FEA results

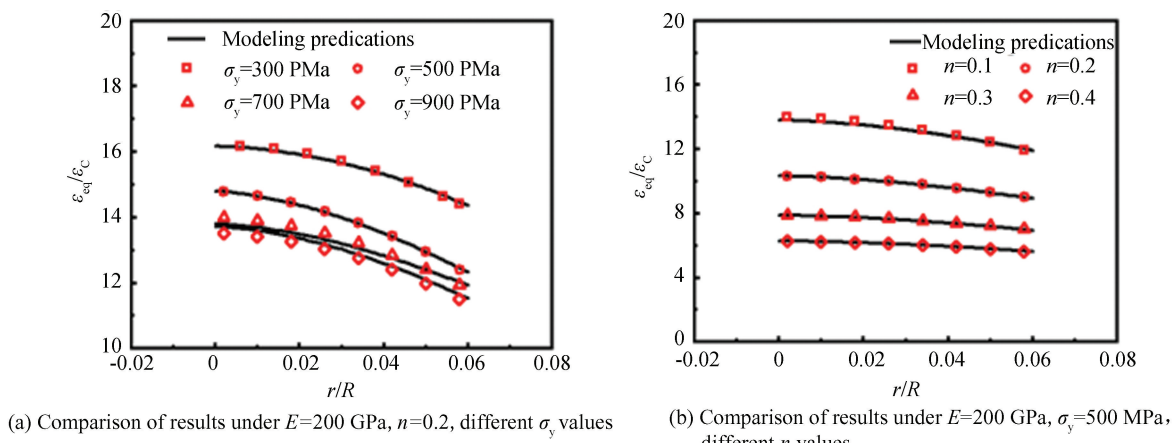


Fig.9 Comparison of strain field model prediction results and FEA results

#### 4.1.3 Verification of maximum stress and strain model

The elastic modulus  $E$  of the compressed material was again set at 200 GPa, with yield stress  $\sigma_y$  values ranging from 300 MPa to 900 MPa in increments of 200 MPa, and strain hardening index  $n$  varying from 0.1 to 0.4 in increments of 0.1. The loading displacement was set between 0.06 mm and 0.2 mm, with increments of 0.02 mm. Scatter plots were generated for maximum stress  $\sigma_{eq-max}$  and maximum strain  $\varepsilon_{eq-max}$  about the dimensionless displacement  $h/h^*$ . As described in Section 1.1, the stress and strain values peak at the clamping end, the FEA results are shown in Fig.10.

A comparative analysis between the analytical

predictions derived from Eq. (17) and FEA results are presented in Fig. 11 and Fig. 12. The model demonstrates strong agreement with FEA simulations, with relative errors of  $\leq 6\%$  for maximum stress predictions and  $\leq 4\%$  for maximum strain estimates.

## 4.2 Experimental Methods and Applications

### 4.2.1 A new SPT method for obtaining cyclic steady-state stress-strain relationship

Based on the theoretical framework relating equivalent stress amplitude to equivalent strain amplitude, the stress-strain relationship of a material in a cyclically stable state can be directly determined from displacement and load amplitude data. The specific experimental procedure is described as follows.

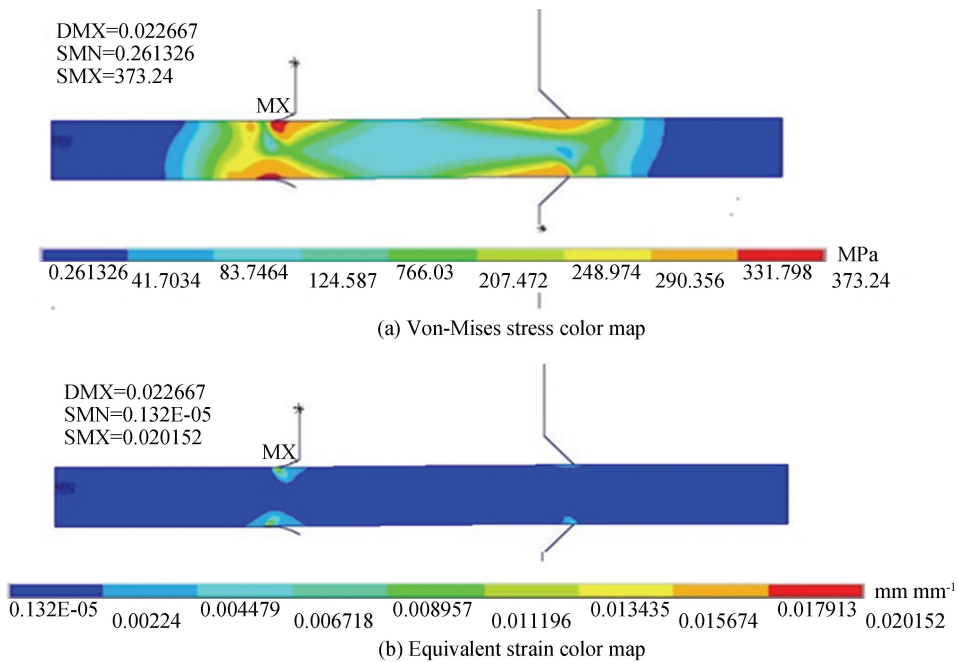


Fig.10 Stress and strain distribution, location and magnitude of Disc<sub>ch</sub> specimen under a loading displacement of 0.06 mm

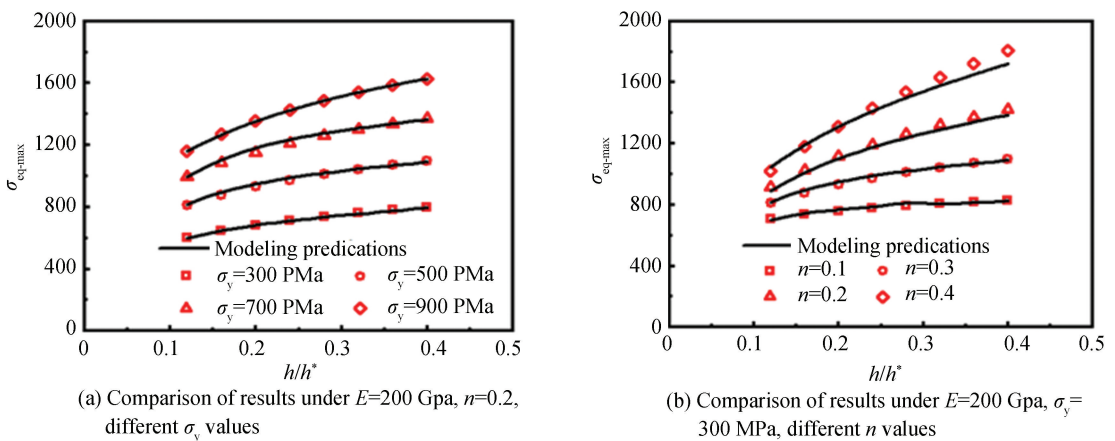


Fig.11 Comparison between the predicted results of the maximum stress model and FEA results

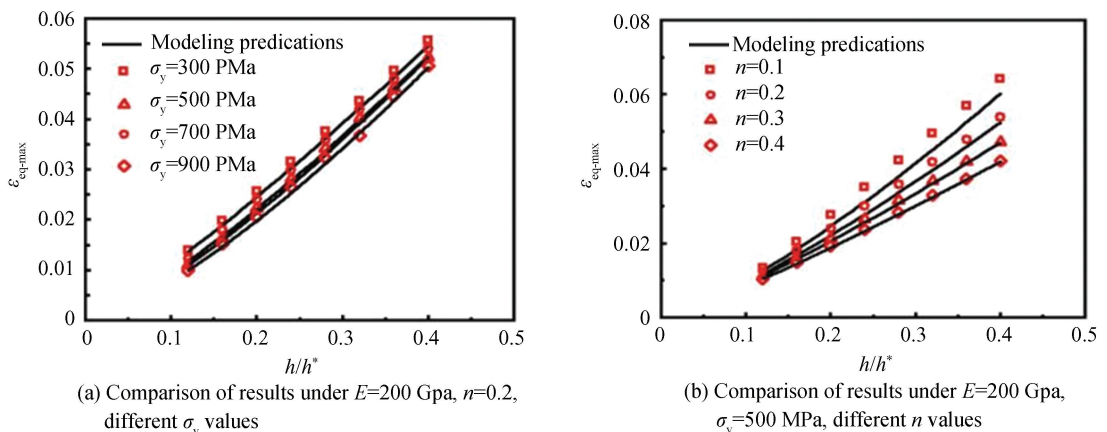


Fig.12 Comparison between the predicted results of the maximum strain model and FEA results

Through experimental data acquisition and processing, stabilized axial displacement-load hysteresis loops were obtained at each controlled displacement amplitude level. The load amplitude ( $P_a$ ) and displacement amplitude ( $h_a$ ) correspond to the maximum cyclic load ( $P_{max}$ ) and maximum displacement ( $h_{max}$ ) measured during testing. By interpolating the vertex coordinates ( $P_{max}, h_{max}$ ) of the hysteresis loops, the displacement-load amplitude curve  $P_a-h_a$  was constructed. The equivalent cyclic stress amplitude  $\sigma_{eq-C-a}$  and equivalent cyclic strain amplitude  $\varepsilon_{eq-C-a}$  were then calculated directly from the  $P_a-h_a$  curve using the theoretical equivalence relationships between stress amplitude and strain amplitude.

For the metal materials 316L and P91, stabilized axial displacement-load hysteresis loops (HL) were obtained from  $Disc_{ch}$  specimens subjected to varying displacement amplitudes. The load amplitude  $P_a$  and displacement amplitude  $h_a$ , derived from the maximum cyclic load  $P_{max}$  and displacement  $h_{max}$ , were used to construct  $P_a-h_a$  curves based on hysteresis loop vertex

data. Fig.13 illustrates the HL and  $P_a-h_a$  curves for  $Disc_{ch}$  specimens made of 316L and P91 materials. The  $\sigma_{eq-C-a}$  and  $\varepsilon_{eq-C-a}$  were derived directly from the  $P_a-h_a$  curve using the theoretical equations for cyclic equivalent stress and strain amplitudes.

Fig.14 presents the cyclic stable  $\sigma_a-\varepsilon_a$  curves and the monotonic  $\sigma-\varepsilon$  curves of the standard round bar specimen, as well as the cyclic stable  $\sigma_{eq-C-a}-\varepsilon_{eq-C-a}$  curves of  $Disc_{ch}$  specimens. The  $\sigma_a-\varepsilon_a$  curves obtained from  $Disc_{ch}$  specimens show strong consistency with experimental data from standard round bar specimens, indicating geometric independence in the steady-state cyclic stress-strain behavior. This agreement confirms that strain-controlled cyclic deformation is governed by intrinsic material properties rather than specimen geometry. Notably, the cyclic  $\sigma_a-\varepsilon_a$  curves for both materials lie above their monotonic counterparts, illustrating the effect of cyclic hardening.

Table 5 summarizes the H-law parameters for both materials obtained using this experimental method under monotonic and cyclic steady-state conditions.

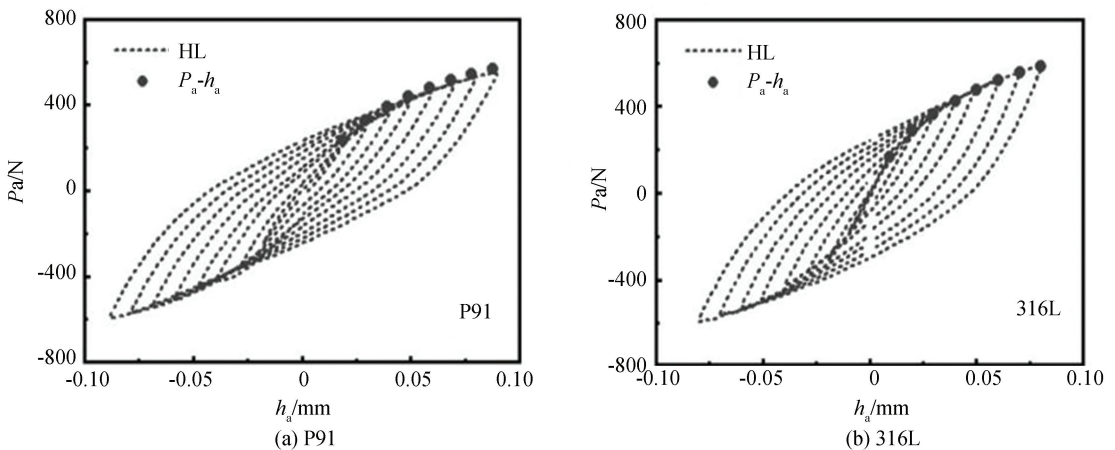


Fig.13  $Disc_{ch}$  sample cyclic stable HL and  $P_a-h_a$  curves

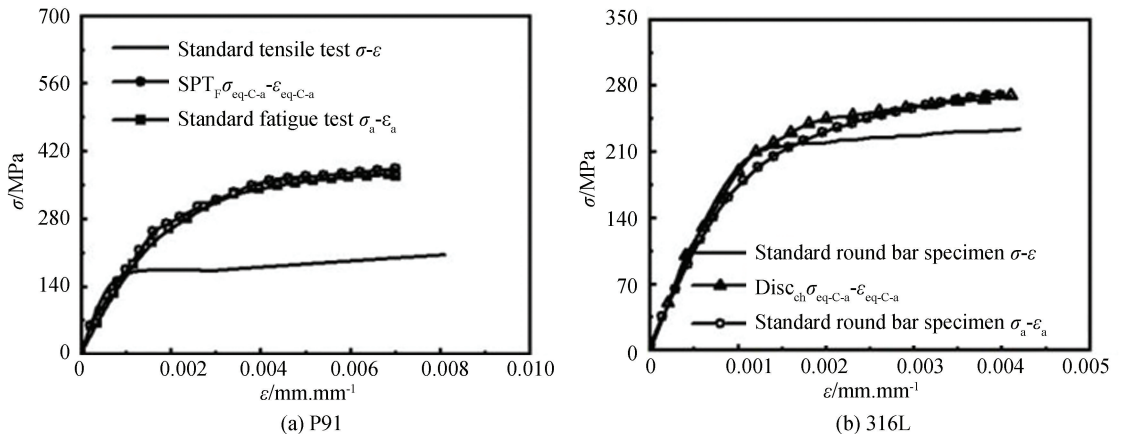


Fig.14 Comparison of cyclic stress-strain amplitude curves between  $Disc_{ch}$  specimens and standard round bar specimens

**Table 5 Monotonic and cyclic constitutive relationship parameters**

Materials	Elastic modulus <i>E</i> /GPa	Monotonic constitutive relationship parameters		Loop constitutive relationship parameters	
		<i>K</i> /MPa	<i>n</i>	<i>K<sub>a</sub></i> /MPa	<i>n<sub>a</sub></i>
P91	224	903.13	0.266	1259	0.237
316L	187	381.15	0.0987	897	0.196

4.2.2 A novel small punch fatigue test methodology

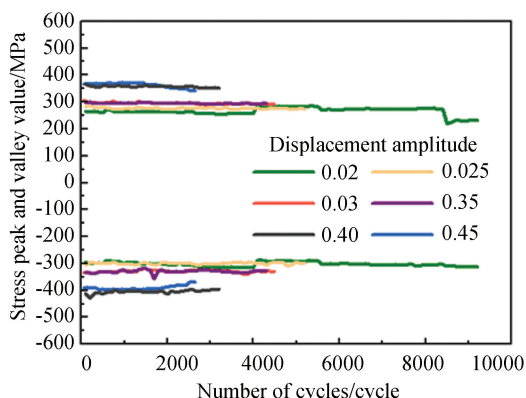
By incorporating the previously determined H-law parameters of the material’s cyclic stress-strain relationship into the maximum stress-strain model, the theoretical stress amplitude and strain amplitude at critical failure locations were calculated for the corresponding load amplitude and loading displacement amplitude. Based on these results, the associated sequences of strain amplitude and displacement-controlled SPT<sub>LCF</sub> life were obtained. Finally, the LCF properties of the material were derived using the Manson-Coffin law. This overall procedure constitutes the novel SPT<sub>LCF</sub> methodology for determining the low-cycle fatigue performance of metallic materials.

Inherent manufacturing tolerances and experimental uncertainties introduce systematic discrepancies between the measured load *P<sub>T</sub>* and theoretical load *P<sub>a</sub>*. Eq.(17) indicates a positive correlation between the load and the strain amplitude. The theoretical strain amplitude  $\epsilon_T$ , stress amplitude  $\sigma_T$ , were empirically calibrated against their experimental counterparts  $\epsilon_a$  and  $\sigma_a$  using the correction formula in Eq.(24),

$$\begin{cases} \epsilon_T = \epsilon_a \frac{P_T}{P_a} \\ \sigma_T = \sigma_a \left(\frac{P_T}{P_a}\right)^n \end{cases} \quad (24)$$

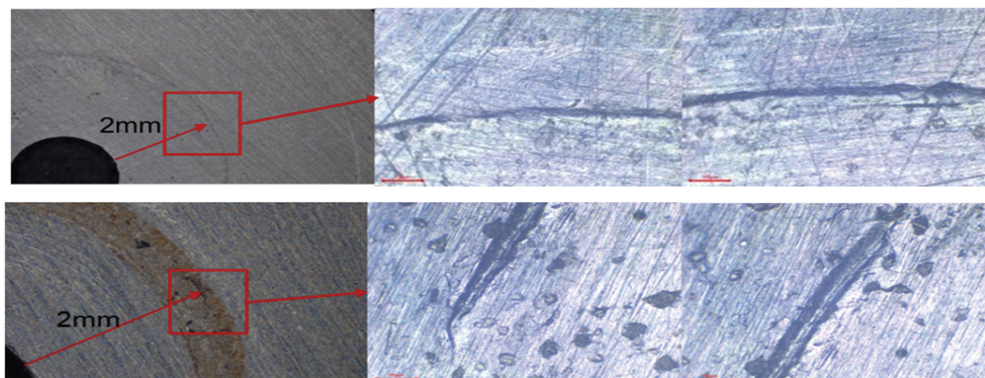
During SPT<sub>LCF</sub>, Disc<sub>ch</sub> specimen was subjected to a multiaxial stress state. Localized crack initiation preferentially occurs at either the upper or lower surface due to stress concentration, leading to progressive material degradation. Nevertheless, due to the symmetrical

loading of Disc<sub>ch</sub>, relative resistance such as frictional force is generated at the crack, with the sample moving up and down. This mechanism causes the overall displacement amplitude to decline more slowly than the peak and valley displacement amplitudes, as illustrated in Fig.15. Therefore, a 5% reduction in the peak/valley displacement amplitude is adopted as the failure criterion for the specimen.



**Fig.15 Diagram of stress amplitude variation with the number of cycles**

Under cyclic loading, localized damage nucleates at the contact endpoints of the flat punch or within the clamping regions of the upper and lower fixtures on the Disc<sub>ch</sub> specimen surfaces. With increasing number of cycles, fatigue cracks propagate progressively, consistent with predictions from the maximum stress-strain model and FEA results. Fig. 16 shows the fractured specimen and the morphology of the localized fatigue cracks.



**Fig.16 Fractured specimen and localized fatigue crack**

Displacement-controlled symmetric cyclic SPT<sub>LCF</sub> were performed on Disc<sub>ch</sub> specimens of P91 and 316L metallic materials. The maximum stress-strain model was employed to calculate stress amplitude and strain amplitude across varying displacement amplitudes, facilitating fatigue life prediction. The relationship between fatigue life  $N_f$  and strain amplitude  $\varepsilon_a$  was quantified using the M-C law, as expressed in Eq.(22).

$$\begin{cases} \varepsilon_{eq-C-a} = \varepsilon_{eq-e-C-a} + \varepsilon_{eq-p-C-a} = \\ \frac{\sigma'_f}{E} (2 N_f)^b + \varepsilon'_f (2 N_f)^c \\ \varepsilon_{eq-e-C-a} = \sigma_{eq-C-a}/E \end{cases} \quad (25)$$

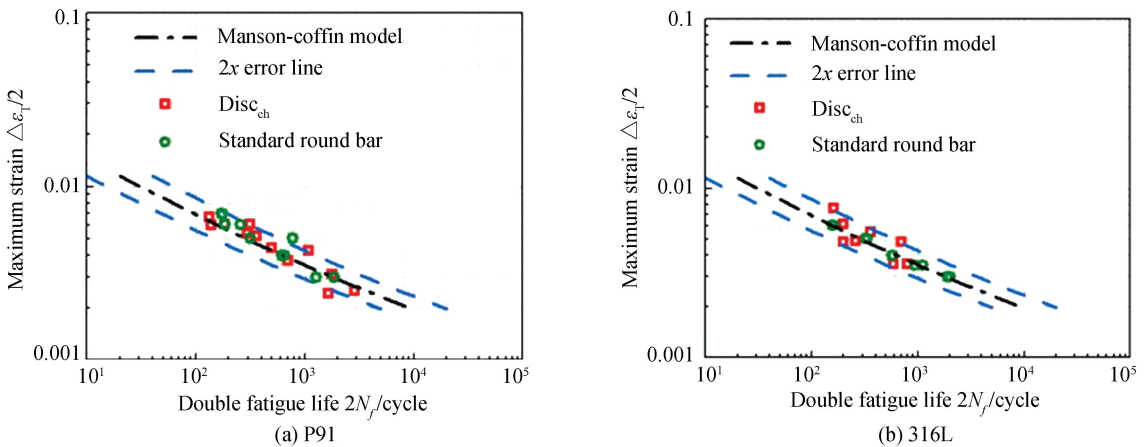
where,  $\varepsilon_{eq-e-C-a} = \sigma_{eq-C-a}/E$  and  $\varepsilon_{eq-p-C-a} = \varepsilon_{eq-C-a} - \varepsilon_{eq-e-C-a}$  are the equivalent elastic strain and plastic strain of the Disc<sub>ch</sub> specimen, where  $N_f$  denotes fatigue life,  $\sigma'_f$  is the fatigue strength coefficient,  $b$  is the fatigue strength exponent,  $\varepsilon'_f$  is the fatigue ductility coefficient, and  $c$  is the fatigue ductility exponent. Table 6 summarizes the M-C law parameter values predicted for Disc<sub>ch</sub> samples. Fig. 17

compares the M-C law fatigue life curves obtained from Disc<sub>ch</sub> specimens and standard round bar specimens, revealing strong alignment. This correlation validates the reliability of the novel testing methodology for LCF performance evaluation.

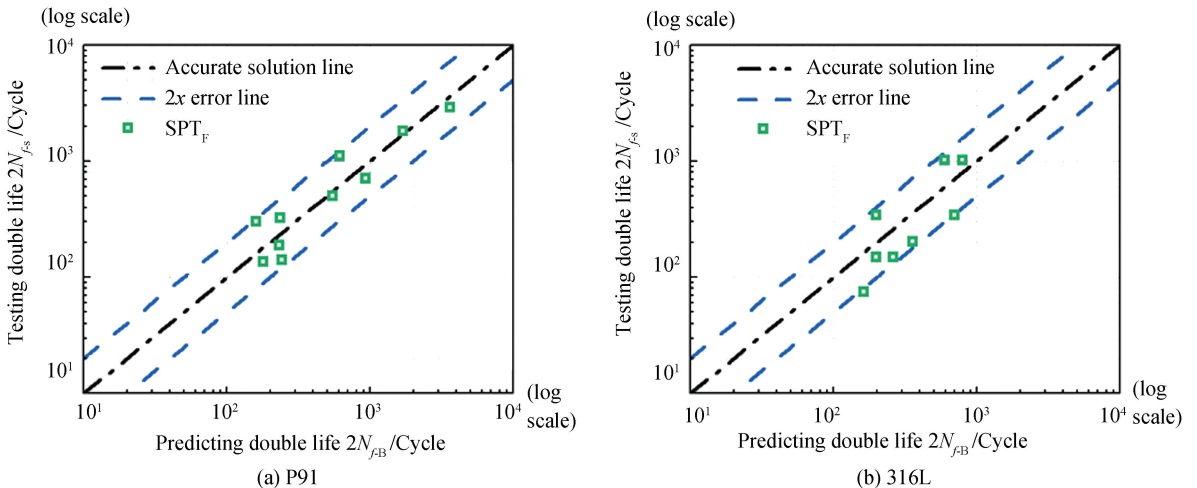
**Table 6** Law parameters of two materials

Material	$\sigma'_f$ /MPa	$b$	$\varepsilon'_f$	$c$
316L	626.4	-0.1079	0.0675	-0.3701
P91	604.8	-0.0658	0.2115	-0.5008

Fig.18 illustrates the relationship between the double life ( $2N_{f-S}$ ) of Disc<sub>ch</sub> specimen (P91 and 316L steels) derived from the novel SPT<sub>LCF</sub> methods and the predicted fatigue life ( $2N_{f-B}$ ) based on standard round bar specimens. The diagonal “exact solution” line represents theoretical equivalence between specimen types, while parallel dashed lines delineate a two-fold safety factor. Experimental data demonstrate that all  $2N_{f-S}$ - $2N_{f-B}$  points fall within the safety boundaries, with deviations complying with acceptable engineering tolerances.



**Fig.17** Fatigue life prediction curve



**Fig.18** Scatter plot of double life curve of SPT<sub>LCF</sub> with Disc<sub>ch</sub> specimen-predicted double life. Although the number of specimens was insufficient to permit a rigorous statistical analysis of

fatigue life scatter, the data obtained from both the small punch fatigue tests and the standard uniaxial fatigue tests (Fig. 17 and Fig. 18) reveal clear and consistent trends. These results demonstrate that the small punch fatigue testing method can be effectively applied to evaluate material fatigue performance.

## 5 Conclusions

This study establishes a combined theoretical and experimental framework for determining the cyclic steady-state stress-strain relationship and evaluating the low-cycle fatigue (LCF) properties of materials using  $Disc_{ch}$  specimens under displacement-controlled small punch fatigue testing ( $SPT_{LCF}$ ). The main findings are summarized as follows:

(1) Four analytical models have been developed: a displacement-load model, an equivalent stress-strain model, a stress-strain field model, and a maximum stress-strain model. Simulations performed on a range of materials show that the predicted stress-strain results deviate by less than 4% from reference values, confirming the accuracy and universality of the proposed models.

(2) Based on the equivalent stress-strain model, a novel displacement-controlled SPT methodology has been developed to determine the stress-strain relationship of materials under cyclic steady-state conditions. To validate this approach, multi-stage symmetric displacement-controlled cyclic SPT tests were carried out on  $Disc_{ch}$  specimens of 316L stainless steel and P91 steel. For comparative purposes, multi-stage symmetric strain-controlled cyclic tests were also conducted on conventional straight round bar specimens. The cyclic steady-state stress-strain curves obtained from the  $Disc_{ch}$  specimens show close agreement with those derived from the standard specimens, with deviations of less than 4%.

(3) Based on the four proposed theoretical models, a conversion relationship has been established to correlate the cyclic steady-state load amplitude and displacement amplitude with the stress amplitude and strain amplitude at the critical failure point of the specimen. By incorporating the local stress-strain approach and substituting the axial stress-strain terms in the Manson-Coffin law with the equivalent stress-strain values at the failure location, a universally applicable  $SPT_{LCF}$  methodology has been developed for  $Disc_{ch}$  specimens across a variety of materials and

geometries.  $SPT_{LCF}$  were performed on  $Disc_{ch}$  specimens of 316L stainless steel and P91 steel. The Manson-Coffin law parameters obtained using the novel  $SPT_{LCF}$  method show good agreement with those derived from standard round bar tests.

This study provides a preliminary validation of the feasibility of the small punch fatigue testing method for evaluating material fatigue performance. It should be noted that the current conclusions are based on a limited sample size, and thus their statistical confidence is constrained. Future work will involve increasing the number of specimens to comply with the ASTM E739-23 standard guide, with the aim of establishing a more statistically significant fatigue life database and design allowable curves.

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